

# Ultra-short dual-core GaAs photonic crystal fiber splitter filled with nematic liquid crystal

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**Abstract.** A simple nematic liquid crystal-filled dual-core photonic crystal fiber polarization beam splitter is designed and analyzed by the full-vector finite element method. The length of the polarization beam splitter is 13.3390  $\mu\text{m}$  and the maximum extinction ratio (ER) is 143.49 dB. The bandwidth of ERs greater than 10 dB is 200 nm at a wavelength of 1.55  $\mu\text{m}$ . This simple and easy-to-manufacture structure has great potential in miniaturized communication devices. © 2021 Society of Photo-Optical Instrumentation Engineers (SPIE) [DOI: [10.1117/1.OE.60.5.056104](https://doi.org/10.1117/1.OE.60.5.056104)]

**Keywords:** optical crystal fibers; polarization beam splitter; coupling length; extinction ratio.

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## 1 Introduction

A dual-core fiber contains a pair of parallel cores exploiting coupling of the evanescent field between the two cores to carry out specific optical functions. Compared with the conventional optical fiber composed of a core and cladding, the twin-core optical fiber has a special structure rendering it useful not only as an optical transmission medium but also in new optical waveguide devices, such as polarization splitters,<sup>1–3</sup> coherent detectors,<sup>4</sup> wavelength division multiplexers,<sup>5</sup> optical switches, and filters.<sup>6,7</sup> Recently, the photonic crystal fiber (PCF)<sup>8–10</sup> has attracted much attention because of its flexible structure, which provides the possibility of designing dual-core fibers and new fiber polarization beam splitters.

The birefringence effect caused by the asymmetry of the dual-core PCF structure yields different coupling lengths and allows separation of different polarization states. In 2003, Zhang and Yang<sup>11</sup> reported a 1.7-mm-long polarization splitter based on dual-core PCFs with extinction ratio (ER) of 11 dB and bandwidth of 40 nm. Since then, polarization beam splitters based on dual-core PCFs have been developed. In 2016, Yong et al.<sup>12</sup> proposed an octagonal dual-core PCF polarization splitter with a bandwidth of 37 nm and splitting ratio of 81.2 dB at 1550 nm. In 2019, Fahmi et al.<sup>13</sup> designed a dual-core splitter with a coupling length of 132  $\mu\text{m}$  and bandwidth of 301 nm. Recently, Úsuga-Restrepo et al.<sup>14</sup> described an all-fiber circular polarization beam splitter with a length of 24.76 mm, ER of 50 dB, and bandwidth of 36.79 nm (ER > 20 dB).

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In general, these beam splitters have a wide bandwidth, large ER, or short length, but none of them can satisfy the requirement of a wide bandwidth, high ER, and short length at the same time.

Herein, a simple “nematic liquid crystal” (NLC)-filled PCF polarization splitter is designed. To enhance the asymmetry of the structure, two circular pores filled with NLC are introduced to the vertical axis in the center, consequently increasing the coupling length of the beam splitter. The full-vector finite element method<sup>15-18</sup> is implemented to analyze the PCF polarization filter. A polarization beam splitter with a length of 13.339 μm, ER of 143.49 dB, and bandwidth of 200 nm (for ER > 10 dB) is achieved at 1550 nm. In addition to these excellent properties, this polarization beam splitter has a simple structure that is easy to manufacture and a performance that can be further improved by optimizing the diameter of each hole.

## 2 Model and Method

The cross section of the dual-core PCF polarization beam splitter is depicted in Fig. 1 in which the white areas are the air holes with diameters of  $d_0$  and  $d_1$  and air hole spacing of  $\Lambda$ . The two central air holes in the optical fiber are circular and have a diameter of  $d_2$ . Therefore, two cores denoted by  $A$  and  $B$  are formed. The gray area is the base material of the optical fiber, that is, gallium arsenide (GaAs). The refractive index is obtained by the following formula:

$$n_{\text{GaAs}}(\lambda) = \sqrt{A_1 + \frac{B_1\lambda^2}{\lambda^2 - \lambda_1^2} + \frac{D\lambda^2}{\lambda^2 - \lambda_2^2} + \frac{F\lambda^2}{\lambda^2 - \lambda_3^2}}, \quad (1)$$

where  $A_1 = 5.372514$ ,  $B_1 = 5.466742 \mu\text{m}^2$ ,  $\lambda_1 = 0.4431307 \mu\text{m}$ ,  $D = 0.2429960 \mu\text{m}^2$ ,  $\lambda_2 = 0.8746453 \mu\text{m}$ ,  $F = 1.957522 \mu\text{m}^2$ , and  $\lambda_3 = 36.9166 \mu\text{m}$ . At a wavelength of 1.55 μm, the refractive index of GaAs is 2.03. Figure 1 shows the cross section of the splitter with center holes infiltrated with NLC of type E7. The permittivity tensor of the E7 material has the diagonal form  $[n_e^2, n_o^2, n_o^2]$ . The anisotropic NLC materials have two refractive indices, the ordinary index ( $n_o$ ) and the extraordinary index ( $n_e$ ), which are calculated by the following Cauchy formula:<sup>19-21</sup>

$$n_{o(e)}(\lambda) = A_{o(e)} + \frac{B_{o(e)}}{\lambda^2} + \frac{C_{o(e)}}{\lambda^4}, \quad (2)$$

where  $\lambda$  is the incident wavelength and  $A_o$ ,  $B_o$ ,  $C_o$ ,  $A_e$ ,  $B_e$ , and  $C_e$  are the coefficients of the Cauchy model, where  $A_o = 1.6933$ ,  $B_o = 0.0070 \mu\text{m}^2$ ,  $C_o = 0.0004 \mu\text{m}^4$ ,  $A_e = 1.6933$ ,  $B_e = 0.0078 \mu\text{m}^2$ , and  $C_e = 0.0028 \mu\text{m}^4$ .

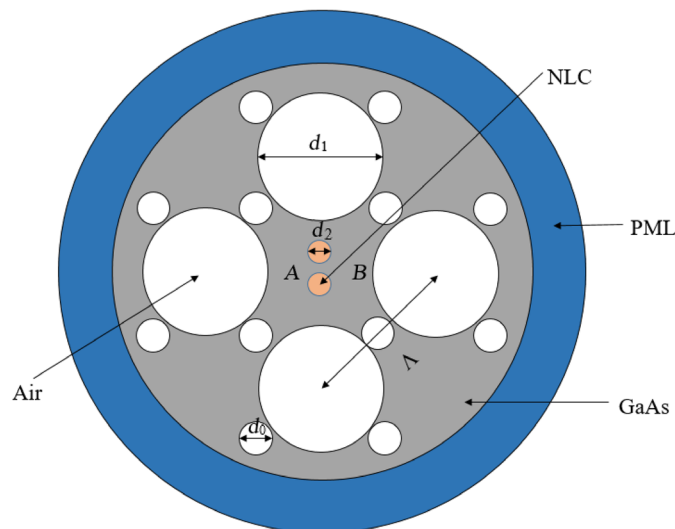


Fig. 1 Cross-sectional structure of the beam splitter.

According to the coupled-mode theory, there are four field distributions in the fiber model, namely the odd and even modes of the  $x$ -polarization and  $y$ -polarization states. The coupling length is<sup>22</sup>

$$L_{x(y)}^\lambda = \frac{\lambda}{2(n_{x(y)}^e - n_{x(y)}^o)}, \quad (3)$$

where  $n_{x(y)}^e$  is the real part of the effective refractive index of the even mode in the  $x$ -polarized and  $y$ -polarized directions, respectively, and  $n_{x(y)}^o$  is the real part of the effective refractive index of the odd modes in the  $x$ -polarized and  $y$ -polarized directions, respectively. When the coupling length of the polarization beam splitter satisfies  $L = L_x/L_y = 0.5$ , the  $x$ -polarization and  $y$ -polarization in the two cores can be separated completely. The confinement loss of the fiber of this mode can be obtained from the imaginary part of the effective refractive index  $n_{\text{eff}}$  of the mode. The relationship between the confinement loss and wavelength of each mode is<sup>23,24</sup>

$$L = 8.686 \times \frac{2\pi}{\lambda} \text{Im}(n_{\text{eff}}) \times 10^6. \quad (4)$$

The relationship between the  $x$ -polarization and  $y$ -polarization output power and input power of the polarization beam splitter is expressed as<sup>25</sup>

$$P_{\text{out}}^{x,y} = P_{\text{in}}^{x,y} \cos^2\left(\frac{\pi}{2} \cdot \frac{L}{L_c^{x,y}}\right), \quad (5)$$

where  $L$  is the transmission distance. The ER is an index that measures the stability of polarization-dependent devices to maintain the polarization state. When a beam of polarized light is incident precisely along a certain optical axis of the polarization-maintaining device, the polarization mode will be excited in the direction of the orthogonal axis. The power ratio of the two orthogonal axes is the extinction ratio of the light source, referred to as the ‘‘ER.’’ The ER is expressed as<sup>26</sup>

$$\text{ER} = 10 \log_{10}\left(\frac{P_{\text{out}}^x}{P_{\text{out}}^y}\right), \quad (6)$$

where  $P_{\text{out}}^x$  and  $P_{\text{out}}^y$  represent the output power of  $x$ -polarization and  $y$ -polarization, respectively.

### 3 Results and Discussion

Figure 2 shows the dispersion curve and loss spectra. As the wavelength increases, the effective refractive indices of the four supermodes decrease and losses increase, although the overall loss is still very small.

While keeping the other parameters unchanged, the relationship between  $\Lambda$  (changing from 1.2 to 1.8  $\mu\text{m}$  in steps of 0.1  $\mu\text{m}$ ) and the coupling length is shown in Fig. 3. The coupling lengths in the  $x$ -polarized and  $y$ -polarized directions increase nonlinearly with air hole pitches ( $\Lambda$ ) at 1.55  $\mu\text{m}$ . As  $\Lambda$  increases, the distance between the two cores increases. At the same time, the coupling length of the polarization beam splitter increases and the coupling effect decreases. However,  $\Lambda$  cannot be reduced indefinitely, which increases the difficulty in production. By considering all of the factors, the optimal value of  $\Lambda$  is 1.2  $\mu\text{m}$ . Similarly, while keeping the other parameters unchanged, when  $d_0$  is changed from 0.21 to 0.26  $\mu\text{m}$  in steps of 0.01  $\mu\text{m}$ , the relationship between the coupling length and coupling length ratio (CLR) is shown in Fig. 4. The coupling lengths in both the  $x$  and  $y$  directions diminish as  $d_0$  increases, and when  $d_0 = 0.251 \mu\text{m}$ , CLR is 0.5.

Figure 5 shows the coupling length in the  $x$ -polarization and  $y$ -polarization directions and CLR for different apertures of  $d_1$ . As the wavelength increases, the coupling lengths for both  $x$  and  $y$  decrease, and when  $d_1 = 1.4 \mu\text{m}$ , CLR is 0.5.

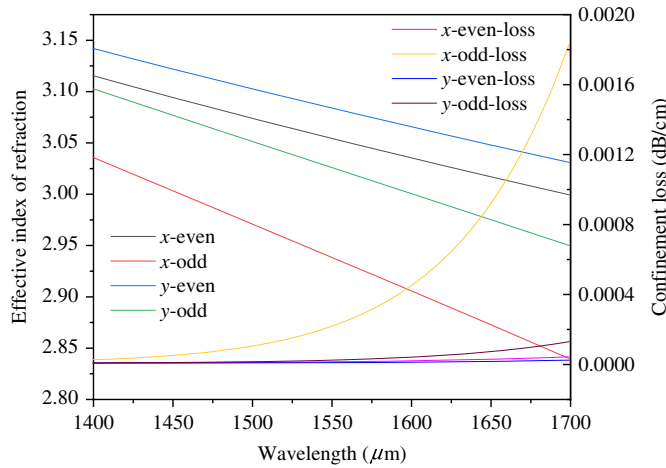


Fig. 2 Dispersion curve and loss spectra.

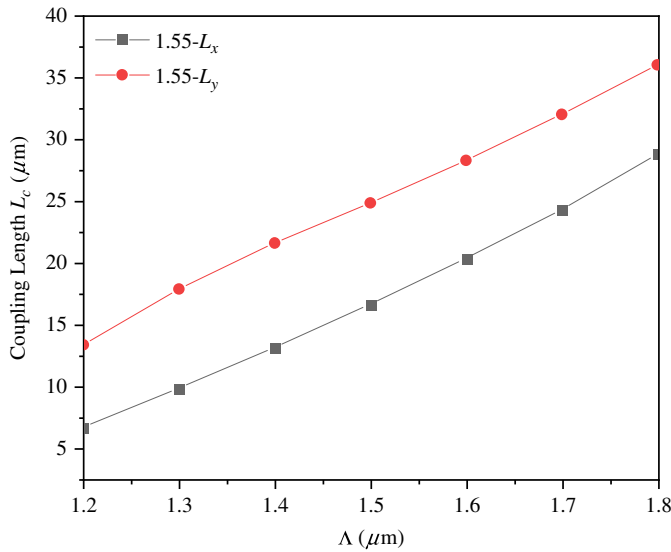


Fig. 3 Influence of  $\Lambda$  on the coupling length.

Figure 6 shows the coupling lengths in the  $x$ -polarization and  $y$ -polarization directions and CLR for different apertures of  $d_2$ . The coupling length in the  $y$ -polarization direction is greater than that in the  $x$ -polarization direction. As the aperture increases, the coupling length increases and coupling between the two cores becomes difficult. The optimal  $d_2$  value is  $0.21 \mu\text{m}$ .

Based on the aforementioned analysis, the optimal parameters of the PCF polarization beam splitter are as follows:  $\Lambda = 1.2 \mu\text{m}$ ,  $d_0 = 0.251 \mu\text{m}$ ,  $d_1 = 1.4 \mu\text{m}$ , and  $d_2 = 0.21 \mu\text{m}$ . The variation of the coupling lengths with wavelength is shown in Fig. 7, and the coupling lengths in the  $x$ - and  $y$ -directions decrease with increasing wavelength. At a wavelength of  $1550 \text{ nm}$ , the coupling lengths in the  $x$ -polarization and  $y$ -polarization directions are  $6.6695 \mu\text{m}$  and  $13.3390 \mu\text{m}$ , respectively. Figure 8 shows the relationship between the normalized power and transmission distance at  $1550 \text{ nm}$ . For a coupling length ratio  $R = L_x/L_y = 0.5$ , the fiber length  $L$  is  $2 \times L_x = 1 \times L_y = 13.3390 \mu\text{m}$  and the transmission distance is  $13.3390 \mu\text{m}$ . That is, the two beams can basically be separated completely.

Figure 9 shows the change of the ERs with wavelengths. The ER is less than  $-10 \text{ dB}$  when the wavelength is  $1.445$  to  $1.645 \mu\text{m}$  and the bandwidth is  $200 \text{ nm}$ . At a wavelength of  $1.55 \mu\text{m}$ , the ER is the largest with the absolute value reaching  $143.49 \text{ dB}$ . To account for possible errors in

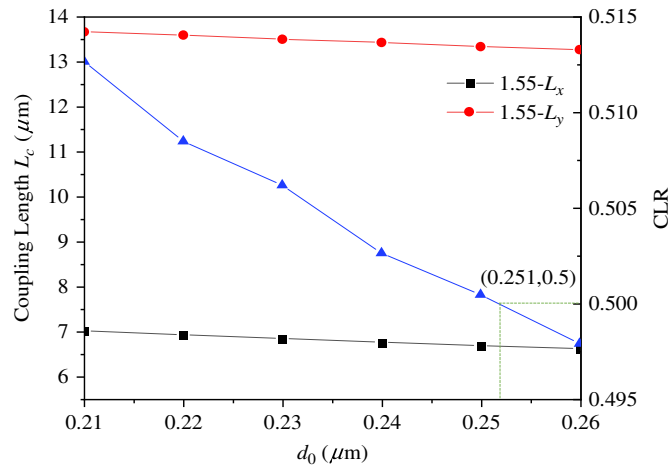


Fig. 4 Influence of  $d_0$  on the coupling length and CLR.

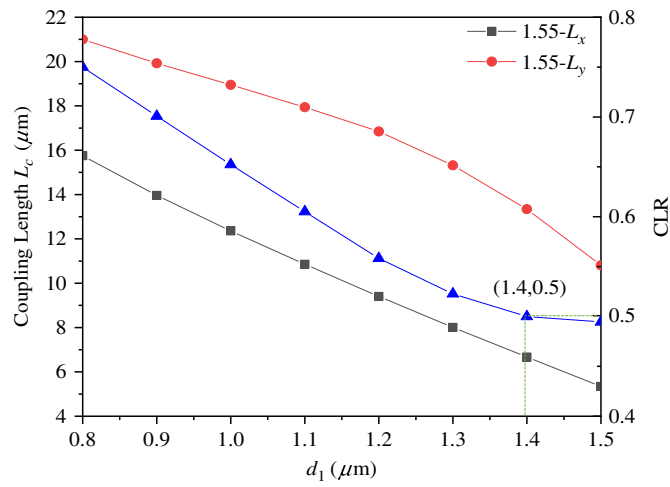


Fig. 5 Influence of  $d_1$  on the coupling length and CLR.

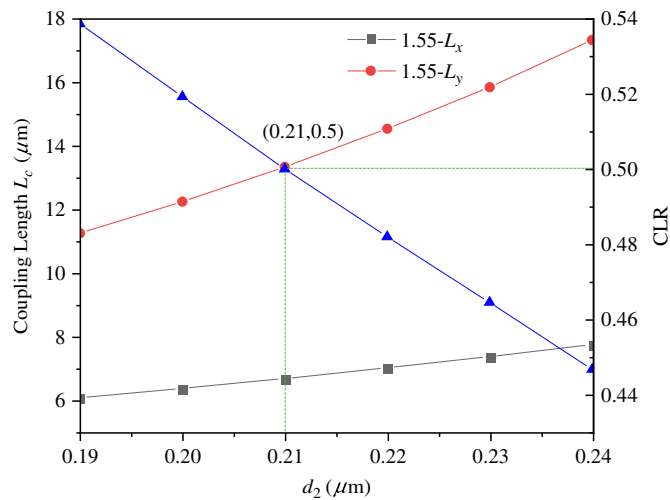
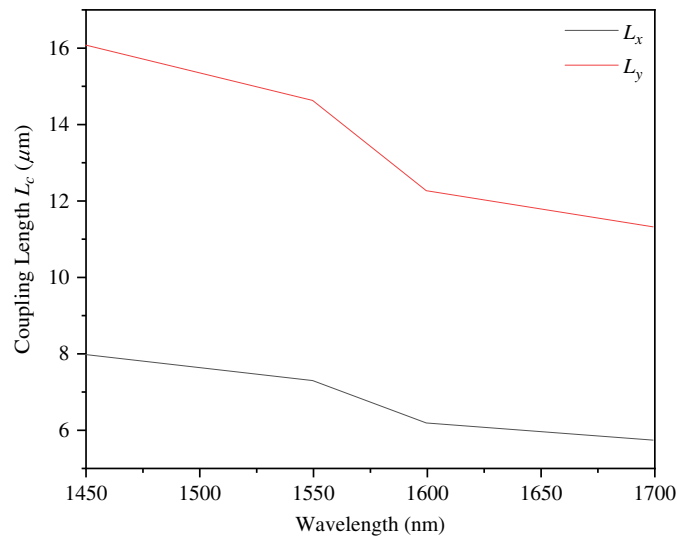
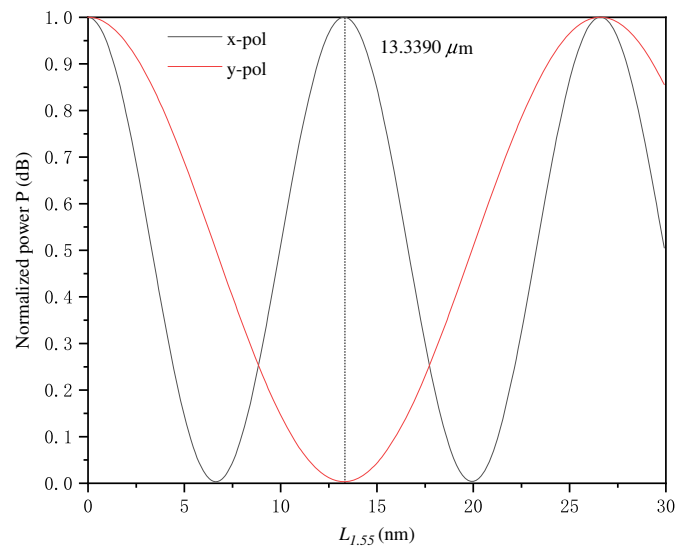


Fig. 6 Influence of  $d_2$  on the coupling length and CLR.



**Fig. 7** Relationship between the coupling lengths of the polarization beam splitter and wavelength under optimal conditions.

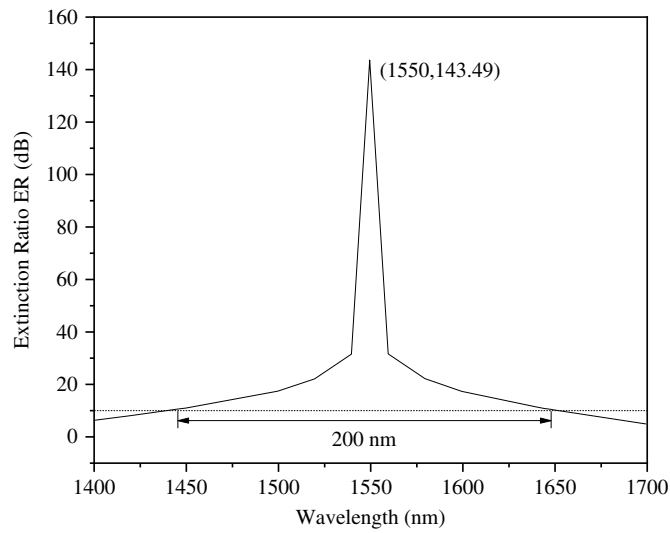


**Fig. 8** Relationship between the normalized power and transmission distance at 1550 nm.

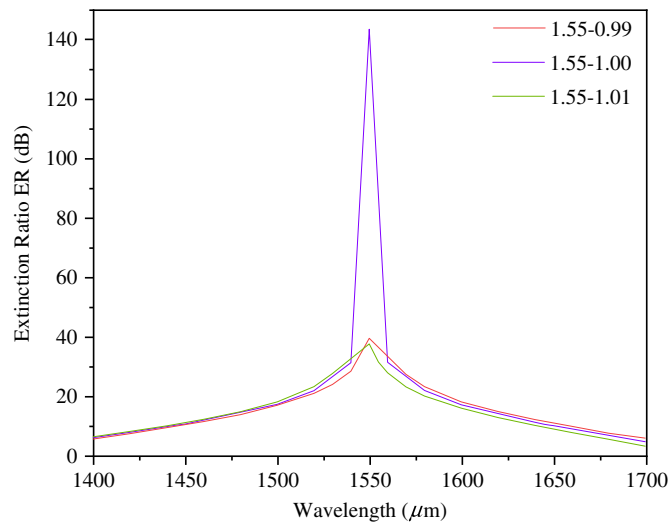
manufacturing, Fig. 10 shows the ER curve with a projected parameter error of  $\pm 1\%$ . An error of  $\pm 1\%$  has almost no effects on the bandwidth, but the maximum ER is affected. All in all, because of the short length of the fiber, the influence of manufacturing errors is minimal.

Figure 11 shows the ER curves of the splitter at different temperatures (20°C, 25°C, and 30°C). By changing the temperature, there is almost no change in bandwidth or ER. Although wavelength and temperature are the two most important factors that affect the refractive index of NLC, the operating temperature can be controlled by a thermoelectric module. GaAs has higher electron mobility and saturation mobility, and it has a unique semi-insulating property. In addition, GaAs materials also have heat-resistant properties, and temperature has little effect on them, giving GaAs materials special uses and versatility; their applications have been extended to areas that silicon dioxide material devices cannot reach.

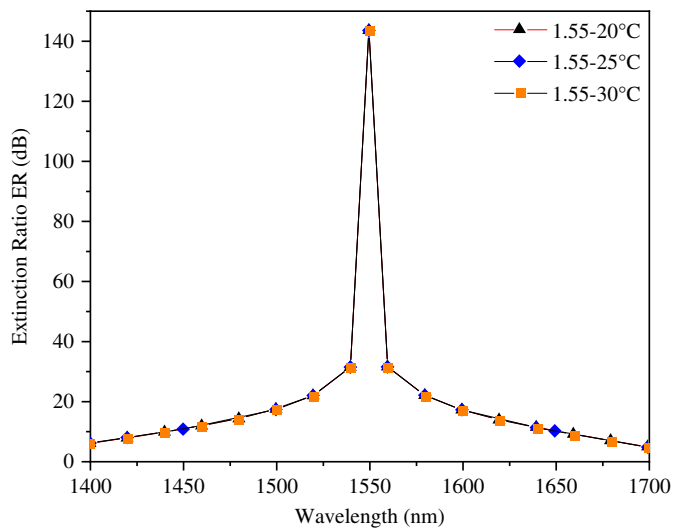
Table 1 compares the performance of this polarization beam splitter with those reported recently in the literature; the virtues of this polarization beam splitter, such as the small size, large ER, and broadband characteristics, are demonstrated.



**Fig. 9** Relationship between ERs and wavelengths.



**Fig. 10** ER curves for parameter error within  $\pm 1\%$ .



**Fig. 11** Influence of different temperatures on ERs of beam splitters at 1550 nm.

**Table 1** Comparison of the properties of the PCF polarization splitter described in this paper with those in the literature.

References	Length of the beam splitter ( $\mu\text{m}$ )	ER (dB)	Bandwidth (nm)
3	732.2	32	15.6
2	401	110.1	140
1	249	-50.7	17
27	89.05	107.21	150
28	47.26	-151.43	104
Our work	13.3390	143.49	200

## 4 Conclusion

A dual-core fiber-filled polarization beam splitter comprising NLC is designed, and the influence of various structural parameters on the coupling characteristics is studied analytically. By changing the diameter of the central and adjacent air holes, a polarizing beam splitter with a short length, ER, and wide bandwidth is obtained. The polarization beam splitter with a length of 13.3390  $\mu\text{m}$  shows a bandwidth of 200 nm with an ER larger than 10 dB, and the largest ER is observed to be 143.49 dB. The performance of the simple and easy-to-manufacture polarization beam splitter is superior to that of similar splitters reported recently in the literature. Further, our results provide the theoretical background to facilitate the design of high-performance crystal fiber polarization beam splitters.

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